

# Origin of Oxide Sensitivity in Gold-Based Catalysts: A First Principle Study of CO Oxidation over Au Supported on Monoclinic and Tetragonal ZrO<sub>2</sub>

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Abstract: The catalytic performance of Au/oxide catalysts can vary significantly upon the change of oxide species or under different catalyst preparation conditions. Due to its complex nature, the physical origin of this phenomenon remains largely unknown. By extensive density functional theory calculations on a model system, CO oxidation on Au/ZrO<sub>2</sub>, this work demonstrates that the oxidation reaction is very sensitive to the oxide structure. The surface structure variation due to the transformation of the oxide phase or the creation of structural defects (e.g., steps) can greatly enhance the activity. We show that CO oxidation on typical Au/ZrO<sub>2</sub> catalysts could be dominated by minority sites, such as monoclinic steps and tetragonal surfaces, the concentration of which is closely related to the size of oxide particle. Importantly, this variation in activity is difficult to understand following the traditional rules based on the O<sub>2</sub> adsorption ability and the oxide reducibility. Instead, electronic structure analyses allow us to rationalize the results and point toward a general measure for  $CO + O_2$  activity, namely the p-bandwidth of  $O_2$ , with important implications for Au/oxide catalysis.

## Introduction

Metal and metal oxide composites constitute some of the most important materials in heterogeneous catalysis.<sup>1-2</sup> The oxide is traditionally regarded as the support merely for dispersing the active metal particles. Owing to the recent findings of the superior catalytic activity of Au/oxide systems,<sup>3-7</sup> the oxide has demonstrated a vital role in catalysis since bare Au is never active enough.<sup>8</sup> There has been considerable interest in the physical origin of the oxide-facilitated Au catalysis.9-17 By performing CO oxidation on clean Au surfaces, <sup>18</sup> Au/TiO<sub>2</sub>, <sup>19-23</sup>

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and Au/MgO,<sup>24,25</sup> recent studies showed that (i) CO adsorbs at the low-coordinated sites of Au particles while O<sub>2</sub> adsorbs strongly at the Au/TiO<sub>2</sub> interface but weakly on pure Au and Au/MgO; (ii) CO oxidation can readily occur at the Au/TiO<sub>2</sub> interface. The O<sub>2</sub> adsorption was proposed to be the key step that influences the catalytic activity of Au/oxide catalysts.<sup>26</sup> This simple framework, however, met difficulties being extended to rationalize a large volume of experimental literature.<sup>27-32</sup> Indeed, it remains unclear as to why the activity depends on the species and structure of the support. Obviously, a clear understanding of this important phenomenon is of huge importance toward designing and optimizing catalysts for emerging technology. Here, our efforts to rationalize the oxide-sensitivity phenomenon

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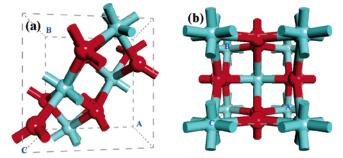
at the atomic level are reported, and the reactant electronic properties are correlated to the activity.

Experimentally, the supporting oxides in Au/oxide catalysts are classified into two types, inert (e.g., MgO, SiO<sub>2</sub>) and active (e.g., TiO<sub>2</sub>, Fe<sub>2</sub>O<sub>3</sub>, IrO<sub>2</sub>).<sup>33,34</sup> The Au/active materials usually exhibit much higher activity than the Au/inert materials for CO oxidation. The active oxides are often reducible, e.g., TiO<sub>2</sub>, while the inert oxides are usually irreducible insulators such as MgO. It was therefore suspected that the reducibility of the oxide makes the difference. However, this suggestion was questioned recently, as new experiments showed that Au supported on  $Al_2O_3$  or  $ZrO_2$  (irreducible oxides) can be quite active for CO oxidation, whereas Au supported on ZnO (reducible oxide) may perform poorly.<sup>29,30</sup> This inconsistency was thought to be caused by different catalyst preparation methods. For example, for Au supported on the irreducible oxide ZrO<sub>2</sub>, a number of experiments have reported a significant variation in activity for CO oxidation.<sup>27–32</sup> Grunwaldt et al. found that the Au/TiO<sub>2</sub> catalyst showed significantly higher activity than the Au/ZrO2 catalyst.28 Using the colloidal deposition methods, Comotti et al. clearly indicated that  $ZrO_2$  is substantially less active than  $TiO_2$ support.<sup>29</sup> However, using laser evaporization, Arrii and coauthors synthesized ZrO<sub>2</sub>- and TiO<sub>2</sub>-supported Au model catalysts. They found that  $TiO_2$  is slightly better than  $ZrO_2$  as support, and the TOF of Au/TiO<sub>2</sub>  $(1.1 \text{ s}^{-1})$  is only 4 times larger than that of Au/ZrO<sub>2</sub> (0.28 s<sup>-1</sup>).<sup>30</sup> Furthermore, the size of the oxide particles was found to be critical to the activity. Zhang et al. showed that Au supported on  $ZrO_2$  nanoparticles (5-15 nm) can yield about 1 order of magnitude higher activity than that of Au supported on the larger ZrO<sub>2</sub> particles (40-200 nm).<sup>31</sup> All these new experimental findings implied that certain, as yet undetermined, fundamental properties of the oxide are vital to activity.

Aiming to build a complete framework for Au/oxide catalysis, in this work CO oxidation over Au supported on various ZrO<sub>2</sub> surfaces (Au/ZrO<sub>2</sub>) is investigated using first principles density functional theory (DFT) calculations. Our results provide a detailed dataset to compare with previous theoretical results on Au/TiO<sub>2</sub> and Au/MgO systems.<sup>19,24</sup> Although the previous surface science studies indicated that oxygen vacancies on the support (such as the F center on MgO) may increase the activity of the Au/oxide catalysts dramatically due to the charging effect of Au,<sup>9,11,12</sup> this work is concerned with nearly neutral Au particles on differently structured oxide surfaces, which should be more relevant to "real catalysts" prepared by typical chemical methods. With several ZrO<sub>2</sub> surfaces being considered as substrates, including two crystal phases and their structural defects, the oxide structure sensitivity in Au/ZrO<sub>2</sub> catalysis is identified and rationalized. Such oxide structure effects could be common to oxide-supported metal catalysts and thus is of general interest to chemistry.

#### **Computational Method and Model**

All DFT slab calculations were performed using the SIESTA package with numerical atomic orbital basis sets and Troullier–Martins norm-conserving pseudopotentials.<sup>35–37</sup> The exchange-correlation functional



*Figure 1.* Crystal structures of (a) the monoclinic  $ZrO_2$  and (b) the tetragonal  $ZrO_2$ . The red and cyan balls represent O and Zr atoms, respectively.

used was the generalized gradient approximation method, known as GGA-PBE.38 A double- $\zeta$  plus polarization basis (DZP) set was employed. The orbital-confining cutoff radii were determined from an energy shift of 0.01 eV. The energy cutoff for the real space grid used to represent the density was set as 150 Ry. To further speed up calculations, the Kohn-Sham equation was solved by an iterative parallel diagonalization method that utilizes the ScaLAPACK subroutine pdsygvx with two-dimensional block cyclicly distributed matrix.<sup>39</sup> The Broyden method was employed for geometry relaxation until the maximal forces on each relaxed atom was less than 0.1 eV/Å. Transition states (TSs) of the catalytic reactions were searched using the constrained minimization method, where the Broyden method was employed to relax all the degrees of freedom except for the constrained reaction coordinate.40,41 The TSs are identified when (i) the forces on the atoms vanish and (ii) the energy is a maximum along the reaction coordinate, but a minimum with respect to all of the other degrees of freedom.

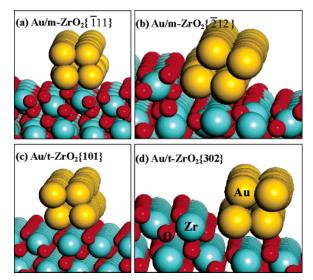
To pinpoint the active site in composite catalysts is a great challenge to both experiment and theory. It is well-known that synthesized ZrO<sub>2</sub> is often a mixture of two phases, monoclinic (m-ZrO<sub>2</sub>, *P*21/*c* symmetry) and tetragonal phases (t-ZrO<sub>2</sub>, P42/nmc symmetry) (see Figure 1).<sup>42</sup> m-ZrO2 is a more stable phase than t-ZrO2, whereas the t-ZrO2 phase is more populated with small zirconia particles.31 Therefore, we anticipated addressing the activity of Au/ZrO2, both phases together with their defects having to be taken into account. Four different low-Miller-index surfaces were selected as substrates, namely, the flat ZrO<sub>2</sub> surfaces m-ZrO<sub>2</sub>{111} and t-ZrO<sub>2</sub>{101}, and the stepped surfaces m-ZrO<sub>2</sub>{212} and t-ZrO<sub>2</sub>{302}. The two flat surfaces are the most stable surface planes in each phase, respectively. m-ZrO<sub>2</sub>{212} contains m-ZrO<sub>2</sub>{101} steps and m-ZrO<sub>2</sub>{111} terraces, while t-ZrO<sub>2</sub>{302} is constituted by t-ZrO<sub>2</sub> {100} steps and t-ZrO<sub>2</sub>{101} terraces. As the stepped m-ZrO<sub>2</sub>{ $\overline{101}$ } and t-ZrO<sub>2</sub>{100} are the second-lowest surface energy planes in each phase, respectively, these types of steps should be the most common structural defects in ZrO<sub>2</sub> particles.<sup>43</sup> Both oxides are insulating. The calculated band gaps are 3.41 eV (exp 4.2 eV) and 3.86 eV (exp 4.2 eV) for m-ZrO<sub>2</sub> and t-ZrO<sub>2</sub> oxide, respectively, which are in good agreement with previous theoretical results with planewave methods.44,45

The Au/ZrO<sub>2</sub> systems were then modeled by adding a two-layer strip of Au on the top of the stoichiometric oxide surfaces (see Figure 2). Similar methods have been successfully applied to study the catalytic

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*Figure 2.* Optimized structures for the Au strip on the four different  $ZrO_2$  surfaces. The red, cyan, and yellow balls represent O, Zr, and Au atoms, respectively.

**Table 1.** Adsorption Energy  $E_{ad}(O_2)$ , the Bond Distance  $d_{O-O}$ , and the Mulliken Charge  $\Delta Q(O_2)$  of  $O_2$  Adsorbed at the Au/ZrO<sub>2</sub> Interfaces; Change of Mulliken Charge of Au Strip  $\Delta Q(Au)$  and the Support ZrO<sub>2</sub>  $\Delta Q(ZrO_2)$  after the Adsorption of  $O_2$  are Also Listed

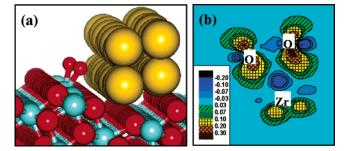
	E <sub>ad</sub> (O <sub>2</sub> )/ eV	d <sub>0-0</sub> / Å	∆ <i>Q</i> (O <sub>2</sub> )/  e	∆ <i>Q</i> (Au)/  e	$\Delta Q(\text{ZrO}_2)/ e $
Au/m-ZrO <sub>2</sub> { $\overline{1}11$ }	0.88	1.475	-0.79	+0.70	+0.09
Au/m-ZrO <sub>2</sub> { 212}	1.07	1.445	-0.68	+0.63	+0.05
Au/t-ZrO <sub>2</sub> {101}	1.08	1.485	-0.79	+0.76	+0.03
Au/t-ZrO <sub>2</sub> {302}	1.71	1.475	-0.77	+0.87	-0.10

activity of Au/TiO<sub>2</sub>.<sup>19</sup> It should be mentioned that the two-layer strip of Au is a reasonable and realistic model based on previous findings: (i) Experiment showed that two-layers of Au particles on TiO<sub>2</sub> are the most effective arrangement for catalyzing CO oxidation.<sup>5</sup> (ii) Theoretically, it was confirmed that appreciable CO and O<sub>2</sub> adsorption is observed only at the low coordinated Au atoms.<sup>18</sup> In the study of CO oxidation on Au/TiO<sub>2</sub>,<sup>19,22</sup> and Au/MgO,<sup>24,25</sup> it was found that the most feasible pathway for the reaction involves CO adsorption at the stepedge of the second-layer Au and O<sub>2</sub> adsorption at the Au/oxide interface. Thus, a two-layer model is sufficient to capture the important chemistry occurring at the interface.

The Au/ZrO2 structures were initially relaxed using Nose thermostat molecular dynamics at 200 K for  $\sim$ 1000 fs, which is essential to locate the best geometry for the Au strip on the substrate. The lattice match between the optimized Au strip and the supporting oxide is generally good, being within  $\sim 1\%$  mismatch. This is mainly due to the large unit cell of substrate surfaces utilized: m-ZrO<sub>2</sub>{ $\overline{111}$ } [p(2 × 2), 13.63  $\text{\AA} \times 14.75 \text{ Å}$ ], t-ZrO<sub>2</sub>{101} [p(3 × 4), 19.19 Å × 14.48 Å], m-ZrO<sub>2</sub>- $\{212\}$  [p(3 × 1), 20.44 Å × 11.69 Å], and t-ZrO<sub>2</sub> $\{302\}$  [p(1 × 4), 17.40 Å  $\times$  14.48 Å]. Because of the large unit cell used and the insulating oxide support, only the  $\Gamma$ -point was used to sample the Brillouin zone in our calculations. The convergence of adsorption energy and reaction barrier with respect to k-point sampling has further been checked by utilizing  $(2 \times 2 \times 1)$  k-point set (the difference is found to be within 0.1 eV). Other calculation details are as those described in the previous work,<sup>19</sup> where the accuracy of the SIESTA method was carefully benchmarked with a plane-wave methodology.

### **Results and Discussion**

Figure 2 shows the optimized structures of the Au strip on the four different  $ZrO_2$  surfaces. The optimized Au structures exhibit close-packed, (111)-like facets. In general, the Au strip



**Figure 3.** (a) Optimized structure, and (b) the charge density difference  $(e/Å^3)$  plot showing the  $Zr-O_2$  bonding plane of the  $O_2$  adsorption on Au/t- $ZrO_2$ {302}.

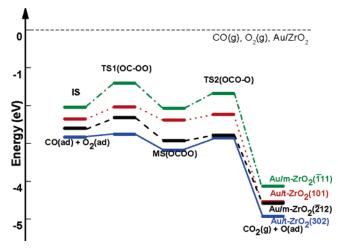
anchors on ZrO<sub>2</sub> through the lattice oxygens. The Au strip adsorption energy is defined as  $E_{ad}(Au) = -(E_{Au/oxide} - E_{oxide} - E_{Au})/N_{first-Au}$ , where  $E_X$  is the DFT total energy of the system X;  $N_{first-Au}$  is the number of the first-layer Au atoms that are in close contact with the oxide. From our calculations,  $E_{ad}(Au)$  is around 0.20 eV for Au/m-ZrO<sub>2</sub>{ $\overline{111}$ }, Au/t-ZrO<sub>2</sub>{101}, and Au/ t-ZrO<sub>2</sub>{302} and 0.26 eV for Au/m-ZrO<sub>2</sub>{ $\overline{212}$ }. This indicates that the Au strips bind weakly on ZrO<sub>2</sub>, in analogy to Au on stoichiometric TiO<sub>2</sub><sup>19</sup> and MgO<sup>46</sup> surfaces. The slightly larger adsorption energy of Au on m-ZrO<sub>2</sub>{ $\overline{212}$ } may be attributed to more two-fold oxygens exposed at step-edges.

Subsequently, O<sub>2</sub> on Au/ZrO<sub>2</sub> was examined by exploring O<sub>2</sub> adsorption on Au, ZrO<sub>2</sub>, and their interfaces. O<sub>2</sub> adsorption energy is defined as  $E_{ad}(O_2) = -[E(O_2/Au/ZrO_2) - E(Au/ZrO_2)]$  $- E(O_2)$ ]). As expected, the most stable adsorption configuration for  $O_2$  is at the Au/ZrO<sub>2</sub> interface. The adsorption energy, the bond distance O-O, and the Mulliken charge of the adsorbed  $O_2$  molecule are listed in Table 1. It can be seen that the  $O_2$ adsorption energy is 0.88 eV at the Au/m-ZrO<sub>2</sub>{111} interface, and it increases to 1.07 and 1.08 eV at the Au/m-ZrO<sub>2</sub>{212} and Au/t-ZrO<sub>2</sub>{101} interfaces, respectively. At the Au/t-ZrO<sub>2</sub>- $\{302\}$  O<sub>2</sub> adsorbs remarkably strongly with an adsorption energy of 1.71 eV (the structure is shown in Figure 3a).<sup>47</sup> The Mulliken charge on the adsorbed  $O_2$  range from -0.68 |e| to -0.79 |e| at the interfaces (Table 1). Because the large electron accumulation destroys the degeneracy of the  $O_2 2\pi^*$  states, the adsorbed O<sub>2</sub> at the interface is spin unpolarized, with a O–O bond length of over 1.40 Å (the gas-phase  $O_2$  distance is 1.24 Å from DFT). This is similar to  $O_2$  adsorption at the Au/rutile-TiO<sub>2</sub>{110} interface, where the adsorption energy is calculated to be 0.86 eV.19

From Mulliken charge analysis (Table 1), it is clear that the electron accumulation on  $O_2$  is mainly due to the electron depletion on Au, while the total charge of the ZrO<sub>2</sub> support is perturbed marginally by the presence of  $O_2$ . This may be better viewed from the calculated charge density difference plot in Figure 3b. The plot represents the charge density change that occurs when the gas-phase  $O_2$  is brought onto the Au/t-ZrO<sub>2</sub>-{302} (O<sub>2</sub> adsorption on other systems is similar and thus not shown). The charge density difference is defined as:  $\rho(O_2/Au/ZrO_2) - \rho(O_2) - \rho(Au/ZrO_2)$ , where  $\rho(O_2/Au/ZrO_2)$ ,  $\rho(O_2)$ , and  $\rho(Au/ZrO_2)$  are the total charge densities of  $O_2/Au/ZrO_2$ , a free  $O_2$  molecule, and Au/ZrO<sub>2</sub> systems without changing their atomic positions. This difference is positive in the regions where

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<sup>(47)</sup> It may be borne in mind that the O<sub>2</sub> adsorption energy on metal systems is generally overestimated by DFT (See refs 50, 51).



*Figure 4.* Energy profiles for CO oxidation over Au on four different ZrO<sub>2</sub> surfaces.

**Table 2.** Reaction Barriers for  $CO + O_2 \rightarrow OCOO$  ( $E_a^1$ ), and  $OCOO \rightarrow CO_2 + O$  ( $E_a^2$ ) Steps at Au/Oxide Interfaces, Respectively

oxide	$m-ZrO_{2}\{ \ \bar{1}11\}$	$m-ZrO_{2}\{ \ \bar{2}12\}$	$t-ZrO_{2}\{101\}$	$t-ZrO_{2}{302}$
$E_{a}^{1}$	0.63	0.28	0.32	0.08
$E_{\rm a}{}^2$	0.39	0.15	0.23	0.33
$W_{p-band}a$	13.47	12.94	12.93	12.14

 $^{\it a}$   $W_{p-band}$  is the calculated  $O_2$  p-bandwidth (see text for definition). The units are eV.

the electronic charge accumulates upon adsorption, and negative where the charge depletes. It is obvious that upon adsorption electrons accumulate on the  $O_2 2\pi^*$  orbitals. The electron density on Zr is highly polarized to maximize the electrostatic interaction between the negative  $O_2$  and the Zr cation.

Interestingly, our results show that the reducibility of the supporting oxide is not necessarily connected to its ability to promote  $O_2$  adsorption. As an irreducible oxide,  $ZrO_2$  performs very well to facilitate  $O_2$  adsorption, comparable to the Au/TiO<sub>2</sub> interface (0.86 eV).<sup>19</sup> By contrast, the irreducible oxide MgO is a poor support since the calculated  $O_2$  adsorption energy at the Au/MgO interface (~0.2 eV)<sup>24,25</sup> is not far from  $O_2$  at the bare Au (~0.0 eV).<sup>18</sup>

We now consider CO oxidation at the Au/ZrO<sub>2</sub> interface. There are apparently two types of mechanisms: (i) CO reacts directly with molecular O<sub>2</sub>, and (ii) O<sub>2</sub> dissociates first and CO then reacts with atomic O. Since our calculated O<sub>2</sub> dissociation barriers at the Au/ZrO<sub>2</sub> interface (O<sub>2</sub>  $\rightarrow$  2O) are generally higher than the corresponding CO + O<sub>2</sub> reaction barriers, the atomic oxygen reaction mechanism, i.e., O<sub>2</sub>  $\rightarrow$  2O; CO + O  $\rightarrow$  CO<sub>2</sub>, is not the lowest-energy pathway and thus not considered here. The reaction profiles for CO oxidation over four different Au/ ZrO<sub>2</sub> interfaces are shown in Figure 4, the reaction barriers are listed in Table 2, and the geometrical structures of the initial state, transition states, and intermediate state are displayed in Figure 5.

As shown, CO initially adsorbs at the edge of the second Au layer in all cases. CO can then react with the interface  $O_2$  through a bimolecular transition state (OC-OO). After the transition state, a metastable OCOO complex is formed.<sup>22,24</sup> The OCOO complex exhibits a coplanar structure with its C end (from CO) attaching to Au and the O end (from  $O_2$ ) sitting on Zr. Its internal O-O bond can then break, leading to the

formation of a CO<sub>2</sub> and an O adatom at the interface. The O adatom can easily be removed by adsorbed CO (the reaction barrier is very low,  $\sim 0.05$  eV).

Among the elementary reaction steps, the  $CO + O_2 \rightarrow OCOO$ step involves two molecules, in which the position of O<sub>2</sub> adsorption and the direction that CO attacks O2 are all related to the detailed substrate structure. It is expected that the energy barriers of the CO +  $O_2 \rightarrow OCOO$  step are more substrate sensitive. This is indeed confirmed by the strong variation in barrier  $E_a^{1}$ , ranging from 0.08 to 0.63 eV. On the other hand, the dissociation of OCOO intermediate is less sensitive to the oxide structure, and the barrier  $E_a^2$  differs no more than 0.2 eV from one surface to another. Therefore, the key to enhance CO oxidation activity is to reduce the reaction barrier  $E_a^{1}$  of the first reaction step. It is also interesting to notice that as the intermediates during CO oxidation become more stable (Figure 4), the difference of  $E_a^1$  and  $E_a^2$  ( $E_a^1 - E_a^2$ ) changes gradually from positive  $(E_a^1 > E_a^2)$  to negative  $(E_a^1 < E_a^2)$ . Consequently, the rate-determining step of CO oxidation is expected to switch from the OCOO formation step to the OCOO decomposition step. This is in fact understandable from the basis of the concept of catalysis, known as Sabatier principle: the optimum activity is the compromise between the rate and strength of chemisorption. If the reaction intermediates are too weakly bonded on the surface, the first step (typically dissociation of reactants) would be too slow to occur; however, if the reaction intermediates are too tightly bonded on the surface, the desorption of products (typically recombination reactions) would be rate limiting. This phenomenon has recently been well addressed by Norskov et al. in ammonia synthesis.48

By comparing the reaction barriers for CO oxidation on the surfaces, one can conclude that the monoclinic stepped surface with the highest barrier being only 0.28 eV is the best catalyst among all the surfaces considered. Our results indicate that CO oxidation on Au/ZrO<sub>2</sub> may be dominated by a small number of stepped sites of monoclinic ZrO<sub>2</sub> considering that the monoclinic phase is the most stable phase of ZrO2. According to the rate equation TOF =  $k e^{(-E_a/RT)}$ [sites], the reaction rate of the CO + O<sub>2</sub> reaction at the stepped sites (room temperature) are at least 3 orders of magnitude larger than that at the flat surfaces, assuming the concentration of stepped sites is only 1% of that of the flat surface sites and the pre-exponential factors k are identical. In addition, the presence of the less stable tetragonal phases can also enhance the activity since the barriers of CO oxidation at Au/t-ZrO2 interfaces are also much lower than those at Au/m-ZrO<sub>2</sub>{111}. With the higher activity at the stepped monoclinic surfaces and at the tetragonal phase, the rate of CO oxidation is expected to increase with the decrease of  $ZrO_2$ particle size since small ZrO<sub>2</sub> particles will expose more structural defects and also contain a higher concentration of tetragonal phases. This supports the experimental findings,<sup>31</sup> in which Zhang et al. suggested that it is the increase of Au/oxide boundary populations which enhances CO oxidation activity.<sup>31</sup> Our results here provide an explanation from the atomic level, emphasizing the critical role of minority sites, i.e., stepped m-ZrO<sub>2</sub> and t-ZrO<sub>2</sub> surfaces.

To shed light on the relationship between oxide structure and activity, we have first calculated the d-projected density of states

<sup>(48)</sup> Bligaard, T.; Norskov, J. K.; Dahl, S.; Matthiesen, J.; Christensen, C. H.; Sehested, J. J. Catal. 2004, 224, 206.

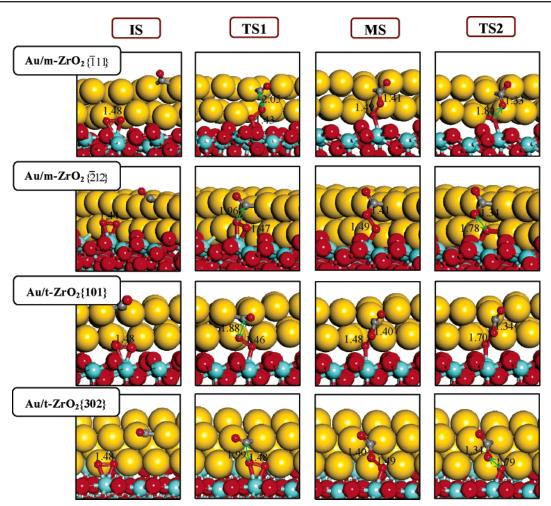


Figure 5. Optimized structures of the initial- (IS), transition- (TS1 and TS2), and intermediate-states (MS) for CO oxidation at four different Au/ZrO<sub>2</sub> interfaces.

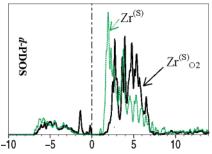


Figure 6. d-Projected density of states of Zr atoms in the  $O_2/Au/t-ZrO_2-$ {302} system. The Fermi level is set as energy zero.

of the surface Zr atoms in the O<sub>2</sub> adsorbed Au/ZrO<sub>2</sub> systems. In each system, the surface Zr bonded with the O<sub>2</sub> (labeled as  $Zr^{(s)}O_2$ ) is selected for analysis to compare with a reference Zr atom that is a normal surface Zr atom (labeled as  $Zr^{(s)}$ ). For illustration, the Zr d-PDOSs of the Au/t-ZrO<sub>2</sub>{302} surface are plotted in Figure 6. It shows that most Zr d states are located in the unoccupied conduction band; due to the extra bonding with O<sub>2</sub>, the d states of  $Zr^{(s)}O_2$  is upshifted in energy compared to the d states of  $Zr^{(s)}O_2$  is upshifted in energy compared to the d states of Zr<sup>(s)</sup>. In the d-PDOS of  $Zr^{(s)}O_2$  there are few extra small peaks around the Fermi level, which are found to be the mixing states between the d-Zr<sup>(s)</sup>O<sub>2</sub> and the O<sub>2</sub>  $2\pi^*$  orbitals. Combining Figure 3b and Figure 6, we can see that the Zr–O<sub>2</sub> interaction through the 4d(Zr) –  $2\pi^*(O_2)$  mixing possesses clear covalent bonding characteristics. Because the d states of Zr are largely empty before and after  $O_2$  adsorption, the electron density difference plot (Figure 3b) does not reveal the rehybridization in the d-orbitals. Nevertheless, the energy change of the Zr d states ( $E_d$ ) can be quantitatively evaluated by

$$E_{\rm d} = \int_{-\infty}^{+\infty} \epsilon \cdot (n_{\rm d}^{O_2 - {\rm ad}} - n_{\rm d}^{\rm bare}) d\epsilon \tag{1}$$

where  $n_d$  are the normalized density of states of Zr (unit: electron/eV) with and without O<sub>2</sub> bonding;  $\epsilon$  is the energy level (eigenvalue). The calculated  $E_d$  is 2.72, 5.88 eV for t-ZrO<sub>2</sub> flat and stepped surfaces, while it is 1.26 and 1.80 eV for m-ZrO<sub>2</sub> flat and stepped surfaces, respectively. For comparison,  $E_d$  for the surface Ti atom in the O<sub>2</sub>/Au/TiO<sub>2</sub> system is also calculated, which is 2.83 eV. As  $E_d$  varies significantly from system to system, it is a qualitative representation of the local activity of metal cation d states. The larger  $E_d$  indicates the more active d states, and vice versa.

From the calculated  $E_d$ , we can see that compared to the monoclinic phase Zr, the tetragonal phase Zr cations can interact more strongly with O<sub>2</sub> and are thus more active. This may not be surprising as t-ZrO<sub>2</sub> is a less stable phase than m-ZrO<sub>2</sub> (the calculated cohesive energy of t-ZrO<sub>2</sub> is 0.11 eV per ZrO<sub>2</sub>-unit less than m-ZrO<sub>2</sub>). Also obvious is that the step-edge Zr cations with generally large  $E_d$  are more active than the flat surface Zr cations. These two trends qualitatively coincide with the

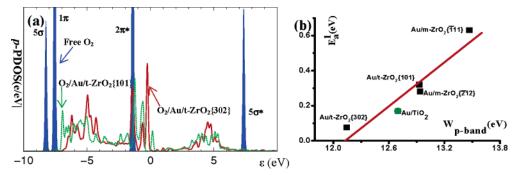


Figure 7. (a) p-Projected density of states for adsorbed O<sub>2</sub> at Au/t-ZrO<sub>2</sub>{302} and Au/t-ZrO<sub>2</sub>{101} (p-PDOS of a free O<sub>2</sub> molecule (spin-unpolarized) is also shown for comparison), and (b) the plot of  $CO + O_2$  reaction barrier ( $E_a^{-1}$ ) against p-bandwidth of  $O_2$ .

determined CO + O<sub>2</sub> activity  $(E_a^1)$  in the Au/ZrO<sub>2</sub> systems. The important role of d states in activating O<sub>2</sub> may also be extended to understand the high activity of TiO<sub>2</sub> and the low activity of MgO. For Au/TiO<sub>2</sub>, the  $E_d$ (Ti) (2.83 eV) of the flat rutile-TiO<sub>2</sub>{110} is quite large, indicating a strong local interaction between O<sub>2</sub> and Ti. In MgO systems, by contrast, where the d states are not available,  $O_2$  is not activated enough, as evidenced by its very low adsorption energy. It should be mentioned that some system-dependent effects such as surface relaxation after O<sub>2</sub> adsorption also contribute significantly to  $E_d$  (e.g., upon O<sub>2</sub> adsorption the existing O-Zr bond will undergo relaxation). This disfavors  $E_d$  as a direct measure for the activity of O<sub>2</sub>.

Alternatively, we have analyzed the p states of adsorbed O<sub>2</sub>, as illustrated in Figure 7a. Similar to its gas-phase counterpart, the p-projected density of states (p-PDOS) of an adsorbed  $O_2$ constitutes three main regions: (a)  $O_2$ -bonding region (mainly  $5\sigma$  and  $1\pi$ ) from -7.5 to -3 eV; (b) O<sub>2</sub>  $2\pi^*$  antibonding region from -3 to +0.5 eV; (c) O<sub>2</sub>  $2\pi^*$  and  $5\sigma^*$  antibonding region: from +0.5 to +7 eV. From the p-PDOSs, two obvious features of adsorbed O<sub>2</sub> can be singled out as compared to that of a free O<sub>2</sub>: first, the broadening of individual O<sub>2</sub> molecular states and, second, the decreased width of the whole p-band. The first feature is universal for molecules interacting with solid surfaces as predicted by the News-Anderson model.<sup>49</sup> The second can also be explained by considering the  $O_2$  bond weakening upon adsorption, which reduces the energy splitting between its bonding (e.g.,  $5\sigma$ ) and antibonding states (e.g.,  $5\sigma^*$ ).

On finding this, we have tentatively plotted the calculated  $CO + O_2$  reaction barrier ( $E_a^1$ ) against the  $O_2$  p-bandwidth  $(W_{p-band})$  in Figure 7b.  $W_{p-band}$  is defined as  $W_{p-band} = \epsilon_{C}^{5\sigma^*} - \epsilon_{C}^{5\sigma^*}$  $\epsilon_{\rm C}^{5\sigma}$ , i.e. the energy difference between the band centers ( $\epsilon_{\rm C}$ ) for  $O_2$  5 $\sigma$  and 5 $\sigma^*$  states. The  $\epsilon_C$  can be approximately calculated using eq 2

$$N_{\rm el} = \int_{-\infty}^{\epsilon_{\rm C}} n(\epsilon) d\epsilon \tag{2}$$

at  $N_{\rm el} = 0.5$  (for  $\epsilon_{\rm C}^{5\sigma}$ ) and  $N_{\rm el} = 11.5$  (for  $\epsilon_{\rm C}^{5\sigma^*}$ ), where  $n(\epsilon)$  is the normalized p-PDOS of O<sub>2</sub> (i.e.,  $N_{\rm el} = 12$  as  $\epsilon \rightarrow +\infty$ ). As shown, the reaction barrier height, including that of Au/TiO2 system ( $W_{p-band}$  for O<sub>2</sub>/Au/TiO<sub>2</sub> is 12.70 eV), can be correlated quite nicely with  $W_{p-band}$ . This implies that it is possible to predict the catalytic activity by evaluating simply the electronic structure of reactants, i.e., the adsorbed O<sub>2</sub>.

It should be pointed out that, in CO oxidation on Au/oxide systems (including ZrO<sub>2</sub>, TiO<sub>2</sub>), as CO adsorbs on the Au particles invariably, it may not then be surprising that the effect of CO is relatively unimportant, which leads to our observation that  $E_a^{1}$  can be simplified as a function of an O<sub>2</sub> electronic property. We also noticed that the calculated O<sub>2</sub> distances at the interfaces are rather similar ( $\sim 1.4$  Å, see Figure 5) and thus apparently show no correlation with the calculated  $E_{a}^{1}$ . This indicates that a subtle change in the electronic structure of O<sub>2</sub> will shift the reaction barrier considerably but has little effect on the adsorption geometry.

Our electronic structure analyses highlight the vital roles of the empty d-states of the metal oxide substrate in activating O<sub>2</sub>. Although the d-states in oxides are largely unoccupied, they are able to mix with the low-lying  $O_2 2\pi^*$  orbitals to further lower the O<sub>2</sub>  $2\pi^*$  energy level, which subsequently leads to the electron flow from the Au strip to the  $O_2$ . The resulted p-bandwidth of  $O_2$  can then be used as a fingerprint of  $O_2$ activity. By varying the oxide species and crystal phases, and by creating surface defects, the d-states of the metal oxide can be modified and consequently change the O<sub>2</sub> adsorption behavior at Au/oxide interfaces. This is therefore a possible reason why CO oxidation over Au-based catalysts is sensitive to the catalyst preparation conditions.

#### Conclusions

In summary, this work demonstrates that the oxide crystal phases and structural defects can significantly modify catalyst activity through reaction kinetics in oxide-supported Au catalysis. The activity of the metal cation empty states varies significantly upon the change of oxide species and structures, which consequently results in the observed oxide sensitivity. We find that the  $O_2$  p-bandwidth can be used as a tool to quantitatively predict the Au/oxide activity, whereas the traditional rules for judging the Au/oxide activity such as the oxide reducibility and O<sub>2</sub> adsorption ability are not applicable.

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